

**The Structure of 5-Anilino-3-oxo-2-phenyl-2,3-dihydro-1*H*-pyrazolo[3,4-*d*]thiazole,  
 $C_{16}H_{12}N_4OS$**

BY DOUGLAS L. SMITH

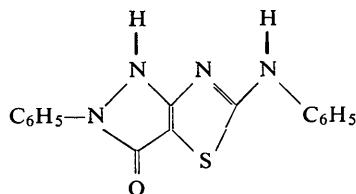
*Research Laboratories, Eastman Kodak Company, Rochester, New York 14650, U.S.A.*

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The crystal structure of  $C_{16}H_{12}N_4OS$  has been determined from three-dimensional X-ray intensity data collected with a punched-card-controlled diffractometer. The structure was solved by the heavy-atom technique and was refined to  $R=3.2\%$  by the method of least squares. Bond lengths show that the molecule has the principal resonance structure, 5-anilino-3-oxo-2-phenyl-2,3-dihydro-1*H*-pyrazolo-[3,4-*d*]thiazole, but there is considerable resonance interaction. The molecules are joined to each other by strong N–H $\cdots$ O and N–H $\cdots$ N hydrogen bonds.

### Introduction

In the course of a study of heterocyclic organic compounds containing sulfur and nitrogen, a compound analyzed as  $C_{16}H_{12}N_4OS$  was prepared. The molecular structure of this compound could not be unambiguously determined by ordinary organic and physical methods. This paper reports the results of the complete solution of the structure of  $C_{16}H_{12}N_4OS$  by X-ray diffraction. The structure ultimately found is 5-anilino-3-oxo-2-phenyl-2,3-dihydro-1*H*-pyrazolo[3,4-*d*]thiazole.



### Crystal data

$C_{16}H_{12}N_4OS$ , M.W. 308.36, m.p. 251–256°C (decomp.). 5-Anilino-3-oxo-2-phenyl-2,3-dihydro-1*H*-pyrazolo-[3,4-*d*]thiazole.

Colorless, thick tabular crystals.

Monoclinic ( $Cu K\alpha_1 = 1.54051 \text{ \AA}$ ),

$$\begin{aligned} a &= 17.107 \pm 0.002 \text{ \AA}, & b &= 7.458 \pm 0.001, \\ c &= 11.285 \pm 0.002; & \beta &= 106.82 \pm 0.01^\circ. \end{aligned}$$

Volume of unit cell,  $1378.1 \pm 0.6 \text{ \AA}^3$ .

Density, calculated ( $Z=4$ ),  $1.486 \text{ g.cm}^{-3}$ ;  
measured (flotation),  $1.50 \text{ g.cm}^{-3}$ .

Linear absorption coefficient,  $\mu = 21.2 \text{ cm}^{-1}$  ( $Cu K\alpha$ ). Total number of electrons per unit cell,  $F(000) = 640$ . Space group,  $P2_1/c$  ( $C_{2h}^5$ ).

Absent spectra:  $h0l$  for  $l$  odd,  $0k0$  for  $k$  odd.

### Experimental

Suitable single crystals of  $C_{16}H_{12}N_4OS$  were supplied by Dr K. R. Rush, of our Laboratories. Space group

extinctions and preliminary lattice dimensions were obtained from precession photographs of a fairly large crystal.

A small rectangular crystal fragment  $0.10 \times 0.10 \times 0.18 \text{ mm}$  was attached to a glass fiber with Canada balsam. The crystal was then used to achieve precise alignment of a Picker four-circle automatic X-ray diffractometer with a Cu-target source and a scintillation detector. Narrow-source, open-counter  $\omega$  scans (Furnas, 1957) through several reflections showed the mosaicity of the crystal to be a uniform  $0.13^\circ$ . After alignment of the diffractometer, 20 reflections from the crystal were centered accurately through a very narrow vertical slit at a take-off angle of  $0.5^\circ$ . Program *PICK2* (Ibers, 1966), modified for the IBM 7044 computer, used these observations as the basis for the least-squares refinement of cell and orientation parameters. These cell parameters, given above, may be affected by systematic errors arising from the manufacture and alignment of the diffractometer and from the shape, perfection, and centering of the crystal, and thus the standard deviations are a measure of precision and not necessarily of accuracy. However, it is thought that their accuracy is sufficiently high to leave the derived interatomic distances unaffected.

Intensity data were collected with Ni-filtered Cu radiation at a take-off angle of  $2.0^\circ$ . The pulse-height analyzer was set for an approximately 90% window. All 2344 unique reflections for  $2\theta < 130^\circ$  were measured by the  $\theta-2\theta$  scan technique (Furnas, 1957) at a  $2\theta$  scan rate of  $0.5^\circ \text{ min}^{-1}$ . The scan range varied from  $1.0^\circ$  at low  $\theta$  to  $1.6^\circ$  at high  $\theta$ . Stationary-crystal stationary-counter background counts of 40 sec were taken at each end of the scan. For count rates above 13000 c.p.s., attenuators were automatically inserted in the diffracted beam by the Picker system.

To check electronic and crystal stability during the period of data collection, the intensity of the 900 reflection was measured every 50th reflection. No systematic drift in this standard was observed. However, the standard deviation of 900, calculated from 56 observations, was 2.0 times what would be expected





on the basis of counting statistics alone. This indicates that some unknown factor is affecting the intensities.

The reduction of the intensity data to structure amplitudes was accomplished with the program *DACOR*, written by the author. The background was approximated by a straight line between the two measured background points. The intensities, corrected for background, were next corrected for Lorentz and polarization effects but not for absorption. Standard deviations  $\sigma(I)$  were based on counting statistics and were corrected to  $\sigma(F)$ . Of the 2344 measured reflections, 313 were less than  $2\sigma(I)$  and were considered unobserved. The intensities of these unobserved reflections were then set equal to  $2\sigma(I)$ . For the 2031 observed reflections, the statistical  $R$  value  $R_s = \sum \sigma(F)/\sum F$  was 0.016.

### Determination and refinement of the structure

The Fourier and least-squares programs used in this analysis were local modifications of programs obtained from Professor J. Trotter (1965). The least-squares program was block-diagonal in coordinates and diagonal in temperature factors. It minimized  $\sum w(|F_o| - |F_c|)^2$ . All least-squares cycles were based on observed reflections only. The atomic scattering factors used for all atoms were those tabulated by Ibers (1962). All computations were performed on the IBM 7044 and 360 computers.

The position of the sulfur atom was determined from an unsharpened three-dimensional Patterson map. A three-dimensional electron density map, phased on the sulfur atom, yielded the positions of the other 21 non-hydrogen atoms. Initially, these 21 atoms were assigned carbon scattering factors. Three cycles of least-squares with individual isotropic thermal parameters and unit weighting reduced  $R_1 = \sum ||F_o| - |F_c|| / \sum |F_o|$  to 0.16

for observed reflections. Inspection of the output temperature factors of cycle 4 enabled the oxygen and the four nitrogen atoms to be identified. This identification was in complete accord with calculated bond lengths and with chemical evidence. Three additional cycles of least-squares reduced  $R_1$  to 0.13. A difference Fourier then indicated considerable anisotropic thermal motion. Several cycles of anisotropic least-squares refinement with Hughes (1941) weighting yielded  $R_1 = 0.062$ . A second difference Fourier clearly showed the 12 hydrogen atoms. The hydrogens were assigned isotropic thermal parameters and refinement was resumed, all parameters being allowed to vary. After several cycles,  $R_1$  converged to 0.032.

In the course of the refinement, between 17 and 26 reflections were considered to be severely enough affected by extinction to be excluded from the refinement.

$$\text{A plot of } (F_c/F_o)^2 \text{ versus } \left( \frac{1 + \cos^4 2\theta}{(1 + \cos^2 2\theta) \sin 2\theta} \right) F_o^2 = f F_o^2$$

yielded the extinction parameter  $g = 2.5 \times 10^{-5}$  (Zachariasen, 1963, 1965). The observed data were then corrected for extinction by the formula  $F_{\text{corr}}^2 = F_o^2(1 + gf/F_o^2)$ . The least-squares analysis was resumed with weights now defined as:

$$w = [\sigma(F_o) + 0.01 F_{\text{corr}} + 0.1(F_{\text{corr}} - F_o)]^{-2}.$$

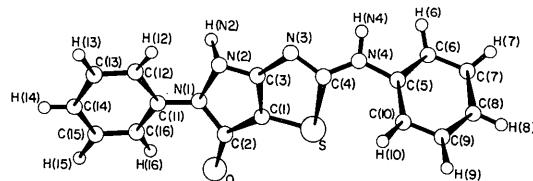


Fig. 1. Configuration and labeling of the molecule.

Table 2. Final parameters of non-hydrogen atoms

$$\text{T.F.} = \exp [-\frac{1}{4}(h^2a^{*2}B_{11} + k^2b^{*2}B_{22} + l^2c^{*2}B_{33} + 2hka^*b^*B_{12} + 2hla^*c^*B_{13} + 2klb^*c^*B_{23})].$$

	$10^5 x$	$10^5 y$	$10^5 z$	$B_{11}$ ( $\text{\AA}^2$ )	$B_{22}$ ( $\text{\AA}^2$ )	$B_{33}$ ( $\text{\AA}^2$ )	$B_{12}$ ( $\text{\AA}^2$ )	$B_{13}$ ( $\text{\AA}^2$ )	$B_{23}$ ( $\text{\AA}^2$ )
S	6842 (2)	38324 (7)	35202 (4)	2.26 (1)	4.65 (2)	1.98 (1)	0.23 (1)	0.81 (1)	0.27 (1)
O	27968 (7)	30099 (20)	51801 (10)	2.86 (4)	6.32 (7)	2.02 (4)	0.35 (5)	0.81 (3)	0.70 (5)
N(1)	29054 (8)	33955 (21)	31868 (11)	2.23 (4)	4.41 (7)	1.80 (5)	0.25 (5)	0.56 (4)	0.26 (5)
N(2)	23952 (8)	39419 (21)	20158 (11)	2.33 (4)	4.53 (7)	1.75 (5)	0.26 (5)	0.47 (4)	0.18 (5)
N(3)	8972 (8)	42868 (22)	13234 (12)	2.48 (4)	4.40 (7)	2.13 (5)	0.24 (5)	0.71 (4)	0.16 (5)
N(4)	-4676 (8)	42978 (23)	12763 (12)	2.34 (4)	5.02 (8)	2.32 (5)	0.22 (5)	0.71 (4)	0.31 (5)
C(1)	16619 (9)	37448 (25)	33767 (14)	2.41 (5)	3.99 (8)	2.28 (6)	0.03 (6)	0.97 (4)	0.15 (6)
C(2)	24741 (9)	33252 (24)	40688 (14)	2.53 (5)	3.71 (8)	2.08 (6)	0.03 (6)	0.80 (4)	0.20 (6)
C(3)	16332 (9)	40210 (24)	21756 (14)	2.42 (5)	3.54 (8)	2.17 (6)	0.03 (6)	0.73 (4)	0.01 (5)
C(4)	3296 (9)	41788 (24)	19088 (14)	2.56 (5)	3.35 (8)	2.23 (6)	0.12 (6)	0.68 (5)	0.06 (5)
C(5)	-11664 (9)	41303 (24)	16944 (15)	2.47 (5)	3.30 (8)	2.92 (6)	0.29 (6)	0.87 (5)	0.28 (6)
C(6)	-18810 (11)	36388 (29)	8105 (17)	2.92 (6)	5.41 (10)	3.62 (7)	-0.37 (7)	0.80 (6)	-0.62 (8)
C(7)	-25986 (12)	35044 (35)	11247 (21)	2.98 (8)	7.40 (14)	6.00 (10)	-1.03 (9)	1.25 (7)	-1.10 (10)
C(8)	-26161 (11)	38450 (34)	23216 (19)	3.45 (6)	6.44 (12)	5.84 (9)	-0.03 (8)	2.53 (6)	0.48 (9)
C(9)	-19111 (11)	43476 (31)	31881 (17)	4.09 (8)	5.26 (10)	3.95 (7)	0.93 (8)	2.09 (5)	0.79 (8)
C(10)	-11806 (10)	45043 (29)	28900 (15)	3.12 (6)	4.84 (9)	3.01 (6)	0.39 (7)	1.15 (5)	0.10 (7)
C(11)	37626 (9)	34357 (24)	34065 (14)	2.36 (5)	3.36 (8)	2.24 (6)	0.10 (6)	0.75 (4)	-0.30 (5)
C(12)	41071 (10)	44348 (27)	26428 (15)	2.92 (6)	4.11 (9)	2.82 (6)	0.01 (6)	0.82 (5)	0.41 (6)
C(13)	49450 (10)	45128 (29)	28964 (16)	3.07 (6)	4.55 (9)	3.71 (7)	-0.22 (7)	1.31 (5)	0.36 (7)
C(14)	54471 (10)	36161 (28)	39056 (17)	2.53 (6)	4.45 (9)	3.84 (7)	-0.01 (7)	0.83 (5)	-0.25 (7)
C(15)	51014 (10)	26004 (28)	46456 (17)	2.79 (6)	4.48 (9)	3.13 (7)	0.69 (7)	0.61 (5)	0.19 (7)
C(16)	42616 (10)	24862 (26)	43974 (15)	2.87 (6)	3.87 (8)	2.67 (6)	0.33 (6)	0.92 (5)	0.26 (6)

The final discrepancy factors for all 2031 observed reflections were  $R_1=0.032$  and  $R_2=[\sum w(F_{\text{corr}} - F_c)^2 / \sum wF_{\text{corr}}^2]^{1/2}=0.039$ . The standard deviation of an observation of unit weight was 1.66. Of the 313 reflections too weak to be observed, only 21 had  $F_c > F_{\text{min}}$  and none had  $F_c > 1.7F_{\text{min}}$ . A final difference Fourier contained values between -0.22 and +0.16 e. $\text{\AA}^{-3}$  and thus verified the structure.

The final values of the observed and calculated structure factors are compared in Table 1. The final positional and thermal parameters and their standard deviations are given for non-hydrogen atoms in Table 2 and for hydrogen atoms in Table 3. The average stan-

dard deviations of the positional parameters expressed in  $\text{\AA}$  are 0.0004 for S, 0.0013 for O, 0.0014 for N, 0.0019 for C, and 0.022 for H. These standard deviations were derived from the semi-diagonalized least-squares process and are to be considered optimistic by a factor of two or three.

### Description and discussion of the structure

The general shape and configuration of the molecule can be seen in Fig. 1. Bond lengths and angles, uncorrected for thermal motion, are given in Tables 4 and 5 and Fig. 2. The standard deviations given in Tables 4

Table 3. Final parameters of hydrogen atoms

	$10^4 x$	$10^4 y$	$10^4 z$	$B$
H(N2)	2483 (12)	3251 (28)	1330 (18)	3.3 (5) $\text{\AA}^2$
H(N4)	-562 (13)	4564 (33)	467 (20)	4.3 (6)
H(6)	-1875 (13)	3365 (31)	-42 (19)	4.0 (6)
H(7)	-3111 (14)	3208 (32)	431 (21)	4.8 (6)
H(8)	-3134 (14)	3793 (35)	2534 (21)	5.3 (7)
H(9)	-1933 (13)	4629 (31)	4039 (18)	4.0 (6)
H(10)	-693 (12)	4906 (31)	3504 (18)	3.5 (5)
H(12)	3740 (12)	5101 (30)	1928 (18)	3.4 (5)
H(13)	5179 (13)	5217 (36)	2336 (20)	4.9 (6)
H(14)	6068 (13)	3738 (29)	4096 (19)	3.3 (5)
H(15)	5448 (12)	1975 (30)	5384 (19)	3.4 (5)
H(16)	4023 (12)	1771 (30)	4935 (18)	3.7 (6)

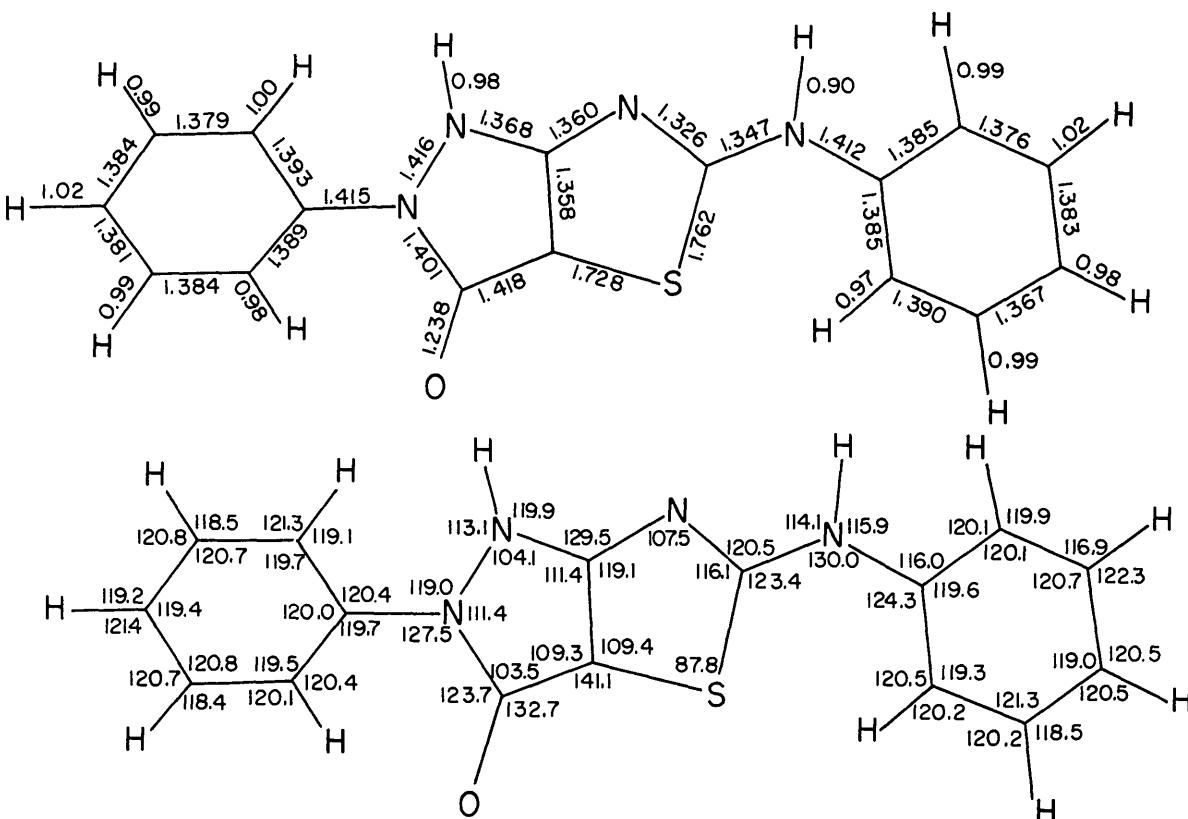


Fig. 2. Bond lengths and angles.

and 5 were calculated from the coordinate standard deviations obtained from the least-squares process. More realistic estimates of the standard deviations were obtained from consideration of the presumably equivalent bond lengths and angles of the two phenyl groups. These standard deviations are 0.007 Å for C-C, 0.017 Å for C-H, 0.7° for C-C-C, and 1.3° for C-C-H.

Table 4. *Intramolecular bond distances*

S—C(1)	1.728 Å	C(11)—C(12)	1.393 Å
S—C(4)	1.762	C(12)—C(13)	1.379
C(1)—C(2)	1.418	C(13)—C(14)	1.384
C(1)—C(3)	1.358	C(14)—C(15)	1.381
C(2)—O	1.238	C(15)—C(16)	1.384
C(2)—N(1)	1.401	C(16)—C(11)	1.389
N(1)—C(11)	1.415	N(2)—H(N2)	0.976
N(1)—N(2)	1.416	N(4)—H(N4)	0.902
N(2)—C(3)	1.368	C(6)—H(6)	0.987
C(3)—N(3)	1.360	C(7)—H(7)	1.017
N(3)—C(4)	1.326	C(8)—H(8)	0.984
C(4)—N(4)	1.347	C(9)—H(9)	0.993
N(4)—C(5)	1.412	C(10)—H(10)	0.965
C(5)—C(6)	1.385	C(12)—H(12)	1.000
C(6)—C(7)	1.376	C(13)—H(13)	0.991
C(7)—C(8)	1.383	C(14)—H(14)	1.024
C(8)—C(9)	1.367	C(15)—H(15)	0.989
C(9)—C(10)	1.390	C(16)—H(16)	0.982
C(10)—C(5)	1.385		

Estimated standard deviations:

Bonds without hydrogen 0.003 Å,

Bonds with hydrogen 0.024 Å.

Inspection of the molecular configuration and the bond lengths reveals that C(1)-C(3) and N(3)-C(4) are double bonds. Therefore, the molecule is 5-anilino-3-oxo-2-phenyl-2,3-dihydro-1*H*-pyrazolo[3,4-*d*]thiazole. However, all bond lengths [except S-C(4)] not involving hydrogen are considerably shorter than normal single-bond distances and it is evident that these bonds have considerable aromatic  $\pi$  character. For example, consideration of the covalent radii for sulfur (Pauling, 1960) and trigonal carbon (Bastiansen & Traetteberg, 1962) leads one to expect the sulfur-carbon single bond to be 1.78 Å and the double bond to be 1.61 Å. Thus the S-C(4) bond (1.762 Å) is not much different from a single bond, but the S-C(1) bond (1.728 Å) has considerable double-bond character and plays an important role in any resonance scheme. Likewise, C(1)-C(3) (1.358 Å) is close to the accepted C=C distance of 1.34 Å, and C(1)-C(2) (1.418 Å) is much shorter than the C-C distance of 1.48 Å characteristic for trigonal carbons (Lide, 1962; Stoicheff, 1962; Bastiansen & Traetteberg, 1962; Sutton, 1965). Carbon-nitrogen bonds are expected to be 1.47 Å for C-N and 1.27 Å for C=N (Sutton, 1965). The N(3)-C(4) bond (1.326 Å) is somewhat longer than a normal double bond, whereas the six other C-N bonds, ranging from 1.347 to 1.415 Å, are considerably shorter than a single bond. The N(1)-N(2) bond (1.416 Å) is somewhat shorter than the 1.45 Å for the N-N single bond in hydrazine (Sutton, 1965) and the C-O bond (1.238) is slightly longer

Table 5. *Intramolecular bond angles*

C(4)—S—C(1)	87.8°	C(16)—C(11)—C(12)	120.0°
S—C(1)—C(3)	109.4	C(11)—C(12)—C(13)	119.7
C(1)—C(3)—N(3)	119.1	C(12)—C(13)—C(14)	120.7
C(3)—N(3)—C(4)	107.5	C(13)—C(14)—C(15)	119.4
N(3)—C(4)—S	116.1	C(14)—C(15)—C(16)	120.8
C(3)—C(1)—C(2)	109.3	C(15)—C(16)—C(11)	119.5
C(1)—C(2)—N(1)	103.5	N(1)—N(2)—H(N2)	113.1
C(2)—N(1)—N(2)	111.4	C(3)—N(2)—H(N2)	119.9
N(1)—N(2)—C(3)	104.1	C(4)—N(4)—H(N4)	114.1
N(2)—C(3)—C(1)	111.4	C(5)—N(4)—H(N4)	115.9
S—C(1)—C(2)	141.1	C(5)—C(6)—H(6)	120.1
N(2)—C(3)—N(3)	129.5	C(7)—C(6)—H(6)	119.9
C(1)—C(2)—O	132.7	C(6)—C(7)—H(7)	116.9
N(1)—C(2)—O	123.7	C(8)—C(7)—H(7)	122.3
N(3)—C(4)—N(4)	120.5	C(7)—C(8)—H(8)	120.5
S—C(4)—N(4)	123.4	C(9)—C(8)—H(8)	120.5
C(4)—N(4)—C(5)	130.0	C(8)—C(9)—H(9)	118.5
N(4)—C(5)—C(6)	116.0	C(10)—C(9)—H(9)	120.2
N(4)—C(5)—C(10)	124.3	C(9)—C(10)—H(10)	120.2
C(10)—C(5)—C(6)	119.6	C(5)—C(10)—H(10)	120.5
C(5)—C(6)—C(7)	120.1	C(11)—C(12)—H(12)	119.1
C(6)—C(7)—C(8)	120.7	C(13)—C(12)—H(12)	121.3
C(7)—C(8)—C(9)	119.0	C(12)—C(13)—H(13)	118.5
C(8)—C(9)—C(10)	121.3	C(14)—C(13)—H(13)	120.8
C(9)—C(10)—C(5)	119.3	C(13)—C(14)—H(14)	119.2
C(2)—N(1)—C(11)	127.5	C(15)—C(14)—H(14)	121.4
N(2)—N(1)—C(11)	119.0	C(14)—C(15)—H(15)	120.7
N(1)—C(11)—C(12)	120.4	C(16)—C(15)—H(15)	118.4
N(1)—C(11)—C(16)	119.7	C(15)—C(16)—H(16)	120.1
		C(11)—C(16)—H(16)	120.4

Estimated standard deviations: Angles without hydrogen 0.2°

Angles with hydrogen 1.3°.

than the 1.22 Å found for C=O in ketones and aldehydes (Sutton, 1965). With these bond-length considerations in mind, it is not surprising to find that atoms C(1), C(2), C(3), C(4), N(1), N(2), N(3), S, and O all lie within 0.07 Å of a plane (Schomaker, Waser, Marsh & Bergman, 1959; Blow, 1960) whose equation is

$$0.0753X + 0.9792Y + 0.1887Z = 3.497,$$

referred to the  $a, b, c^*$  axes. These observations are all consistent with the existence of lone pairs of electrons on the N, S, and O atoms and it is reasonable to expect that the transfer of charge from these lone-pair orbitals into a conjugated system involving double bonds greatly increases the stability of the molecule. Many structures closely related to the present one have been determined with quite good agreement in bond features. Among some recent ones are antabuse (Karle, Estlin & Britts, 1967), uric acid (Ringertz, 1966), iso-cytosine (Sharma & McConnell, 1965), alloxan (Singh, 1965), purine (Watson, Sweet & Marsh, 1965),  $C_{17}H_{15}Cl_3N_3S_2$  (Karle & Karle, 1967) and  $(C_9H_9N_3S)_2$  (Allmann, 1967).

The two phenyl groups are quite normal. The average bond lengths are 1.383 Å for C-C and 0.993 Å for C-H, and the average C-C-C and C-C-H bond angles are both 120.0°. The planes of the phenyls are tilted relative to the approximate plane of S, N(1)···N(3), C(1)···C(4) by 27.8° for the C(5)···C(10) phenyl and by -28.7° for the C(11)···C(16) phenyl. These tilts coupled with distorted angular distributions at N(1), N(4), and C(5) serve to relieve the strain caused by several close intramolecular contacts. The short contacts are C(16)···O (2.914 Å), H(16)···O (2.380 Å), S···C(10) (3.102 Å) and S···H(10) (2.483 Å), all of which are considerably shorter than the sums of the van der Waals radii. The short contacts possibly suggest weak intramolecular hydrogen bonds, but more likely represent a compromise among the requirements of molecular packing, atomic size, and resonance structures.

Fig. 3 shows the molecular packing and hydrogen-bond network in the unit cell, and Table 6 gives the relevant distances and angles. Table 7 lists other inter-

molecular approaches (excluding hydrogen) less than 3.5 Å. Molecules related by the  $c$  glide plane are connected by strong N-H···O bonds to form infinite ribbons. In addition, alternate pairs of N-H···N bonds, in the form of eight-membered rings about centers of symmetry, join each ribbon with the ribbon above and the ribbon below to form sheets of hydrogen-bonded molecules.

Table 6. Hydrogen bond distances and angles

$N(2)_I \cdots O_{II}$	2.776 Å
$H(N2)_I \cdots O_{II}$	1.806
$N(4)_I \cdots N(3)_{III}$	3.002
$H(N4)_I \cdots N(3)_{III}$	2.115
$N(2)_I \cdots H(N2)_I \cdots O_{II}$	171.9°
$H(N2)_I \cdots O_{II} \cdots C(2)_{II}$	133.6
$N(4)_I \cdots H(N4)_I \cdots N(3)_{III}$	167.3
$H(N4)_I \cdots N(3)_{III} \cdots C(4)_{III}$	117.7

I, II and III refer to the equivalent positions  $x, y, z$ ;  $x, \frac{1}{2} - y, -\frac{1}{2} + z$ , and  $-x, 1 - y, -z$ , respectively.

Table 7. Short intermolecular contacts

$C(9)_I \cdots O_{II}$	3.344 Å
$O_I \cdots C(9)_{II}$	3.344
$N(4)_I \cdots S_{III}$	3.417
$C(14)_I \cdots C(16)_{IV}$	3.437
$C(16)_I \cdots C(14)_{IV}$	3.437
$C(10)_I \cdots C(3)_{III}$	3.452
$C(8)_I \cdots N(1)_{III}$	3.453
$C(10)_I \cdots C(1)_{III}$	3.470

I, II, III and IV refer to the equivalent positions  $x, y, z$ ;  $-x, 1 - y, 1 - z$ ;  $-x, \frac{1}{2} + y, \frac{1}{2} - z$ ; and  $1 - x, 1 - y, 1 - z$ , respectively.

In the N-H···O bond, the N···O separation is 2.776 Å, considerably shorter than the normal N-H···O separation of 2.93 Å, given by Wallwork (1962) and somewhat shorter than the normal  $N^+ \cdots H \cdots O^-$  distance of 2.83 Å. This indicates that an important resonance form of the molecule involves species with  $N(2)^+$  and  $O^-$ . In the N-H···N bonds, the N···N separation of 3.002 Å is shorter than the normal 3.07 Å.

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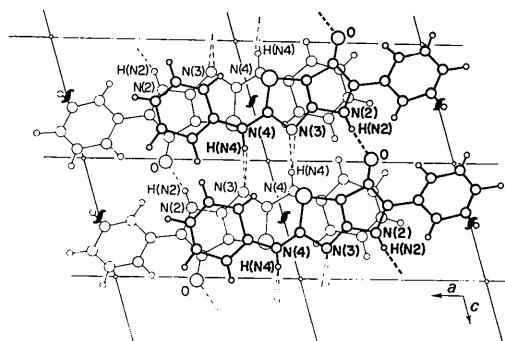


Fig. 3. A view down the  $b$  axis showing contents of a unit cell. Broken lines represent hydrogen bonds.

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## Polytypes of the Two-Packet Chlorites

BY V. A. DRITS AND YU. V. KARAVAN

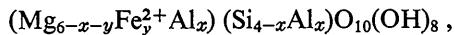
Geological Institute of the Academy of Sciences of the USSR, Pyzhevsky 7, Moscow, USSR

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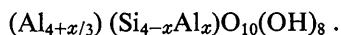
The systematic derivation of all the possible chlorite polytypes with a two-packet unit cell is carried out. The derivation is completed for trioctahedral, dioctahedral and di-trioctahedral chlorites. The diffraction features which should identify the different chlorite polytypes are calculated.

The chlorites are layer-lattice silicates of variable composition. They consist of regularly alternating talc- and brucite-type layers (Pauling, 1930). The talc-type layer is three-storied, consisting of an octahedral net occupying the centre and two tetrahedral nets adjoining the octahedral net from above and below. The brucite-type layer is one-storied, consisting of one octahedral network. Each pair of three-storied and one-storied layers forms a chlorite packet. It would be accepted that a chlorite packet comprises one three-storied layer and two halves of two different one-storied layers, one covering it and one underlying it.

Chlorites are divided into trioctahedral, where all the octahedra are occupied by cations, and dioctahedral, in which only  $\frac{2}{3}$  of all the octahedral positions are occupied. The ideal structural formula for trioctahedral chlorites is



and for dioctahedral chlorites,



The structure of the chlorites is made up from chlorite packets superimposed one upon another in a varying manner. Scientific publications usually describe one-packet and semi-random chlorites (McMurphy, 1934; Zvyagin, 1964; Brown & Bailey, 1962; 1963; Shirozu & Bailey, 1965).

However, McMurphy (1934), and later Garrido (1949) and Brindley, Oughton & Robinson (1950) pointed out the occurrence of chlorites with a different periodicity of packets along the *c* axis, and discussed some polytypes possible for chlorites with more than one packet per repeat period. Chlorites with a two-packet monoclinic unit cell have been described by Drits & Lazarenko (1967) and Wlasow & Drits (1967). This paper gives a systematic determination of all the possible chlorite polytypes with a repeat period of two chlorite packets and their diffraction features are calculated in order to determine the structure of the new two-packet chlorites which we have investigated (Wlasow & Drits, 1967; Drits & Lazarenko, 1967). In deriving these polytypes the analytical method of Zvyagin (1964) was used, providing a simpler solution of this problem.

Table 1. Components of the displacements  $\sigma$  and  $\tau$  along the *a* and *b* axes which are possible in structures of layer silicates

Symbols of the displacements	$\sigma_1$ $\tau_4$	$\sigma_2$ $\tau_5$	$\sigma_3$ $\tau_6$	$\sigma_4$ $\tau_1$	$\sigma_5$ $\tau_2$	$\sigma_6$ $\tau_3$	$\tau_0$	$\tau_+$	$\tau_-$
Components of the displacements	$\frac{1}{3}, \frac{1}{3}$	$-\frac{1}{3}, \frac{1}{3}$	$\frac{1}{3}, 0$	$-\frac{1}{3}, -\frac{1}{3}$	$\frac{1}{3}, -\frac{1}{3}$	$-\frac{1}{3}, 0$	$0, 0$	$0, \frac{1}{3}$	$0, -\frac{1}{3}$

Displacement components are expressed in multiples of *a* and *b*.